

DESIGN OF WAVEGUIDE BANDPASS FILTER IN THE X-FREQUENCY BAND

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How to design a broad bandpass filter in the X-frequency range? The design procedure is developed from synthesis to electromagnetic simulation. The filter to be designed is implemented in waveguide technology with inductive irises coupling sections. So as to evaluate equivalent susceptance of iris aperture, two methods were used: the first one is based on classical mode matching techniques, whereas the second uses electromagnetic-based equivalent circuit. The susceptance values ensued from these methods were then compared. Validations for both methods were carried out using Ansoft-HFSS electromagnetic simulations and measurements.

I. INTRODUCTION

The design of passive microwave functions, and especially filters, is a study domain of the most important interest. Indeed, with the widespread applications of communication systems, the bandpass filter is an essential component. Thus, the quality of filter is extremely important. The development of bandpass filters has emphasized high electrical performances in terms of insertion losses and selectivity. Within this context, designers are faced to several problems linked to the design control, in keeping with the modeling accuracy. The design of such filters requires specific techniques which cover a wide spectrum of knowledge. Students must be introduced to that necessary knowledge that will make them immediately productive upon graduation.

The present paper deals with filter design techniques taught to students in second year of engineering school. During their courses, students get acquainted with filter synthesis techniques, microwave transmission lines analyses, electromagnetic model techniques as well as simulation software, either circuit (Agilent-ADS) or electromagnetic (Ansoft-HFSS) frameworks. The present study recount practical works that gather this entire knowledge and allow students their application in concrete case.

Development of filter design rests on three steps. Application of Tchebycheff theory is first performed. So as to meet the filter specifications, the filter order is determined as well as its lumped element circuit prototype. In expectation of the filter implementation into waveguide technology, that limits the achievable impedance range, synthesis techniques based on impedance inverters and series resonators are then introduced.

Implementation into waveguide technology requires equivalent network for impedance inverters. From circuit analyze point of view, impedance inverter are modeled as inductance set in parallel. Implementation of the whole cascaded elements of the filter is then performed on Matlab. The circuit analyze results evidenced the difference between lumped- (Tchebycheff theory) and semi-lumped- (Richards transformation) elements representation. Hence, so as to meet the expected specifications, further analyses were carried out to define the appropriate filter order.

In a second time, inductances are realized by vertical iris apertures which dimensions are calculated using two methods. The first one is a conventional mode matching technique [1] whereas the second uses analytical approximation based on electromagnetic equivalent circuit [2]-[3], which concepts are detailed in part IV. Validation of both methods is achieved through electromagnetic simulations using Ansoft-HFSS electromagnetic simulator as well as experimental results presented. This approach shows the strong link between lumped circuit and the electromagnetic modelisation which leads suited physical dimensions for the device.

II. BASE-FILTER DESIGN PROCEDURE

A. Lumped Elements Design

The filter to be designed is a Tchebycheff bandpass filter in the X frequency range. Hence, the waveguide used is a WR90 with 0.9 inch x 0.4 inch section. The central frequency of the filter was determined with respect to single mode propagation. Hence, according to the waveguide dimensions, the first propagating mode is a TE_{10} mode with 6.56-GHz cutoff frequency, whereas the first higher-order mode is a TE_{20} with 13.12-GHz cutoff frequency. Hence a central frequency (f_0) of 9.84-GHz was chosen which corresponds to the center frequency of the mono-mode band. Finally, the 40-dB rejection level (L_{ar}) is calculated as in (1) [4]-[5]:

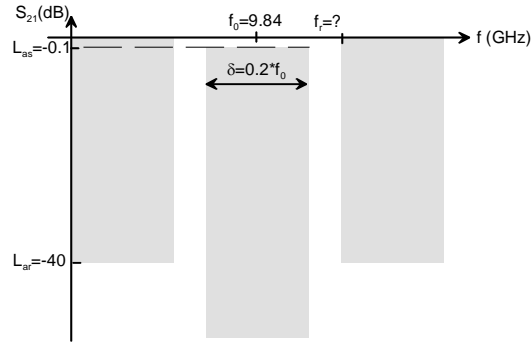


Fig. 1. Bandpass filter specifications.

$$f_r = 1.2 * \left(1 + \frac{\delta}{2}\right) f_0 \quad (1)$$

where δ is the filter relative bandwidth ($\delta=0.2$). The filter specifications are depicted in Fig. 1.

The first conception step consists in the determination of the order of the equivalent low-pass filter prototype, which is the base-network for filter synthesis, whatever the filter type is, high-pass, band-pass or stop-band. It is well known that lowpass prototype filter can be transformed into bandpass response using the frequency substitution (2):

$$\omega' \leftarrow \frac{1}{\delta} * \left(\frac{\omega}{\omega_0} - \frac{\omega_0}{\omega} \right) \quad (2)$$

where δ is the fractional bandwidth of the passband, and ω_0 the center frequency. Application of this transformation with $\omega = \omega_r$, allows to evaluate the reject frequency ($\omega'r$) of the lowpass prototype. Then, the filter order is determined using the equation (2):

$$N \geq ch^{-1} \sqrt{\frac{\frac{Las}{10^{10}} - 1}{\frac{Lar}{10^{10}} - 1}} / ch^{-1}(\omega', r) \quad (3)$$

For the specified criterion a 5th-order filter is necessary. The base lowpass prototype network is depicted in Fig. 2, where normalized inductance and capacitance values assumed to be the Tchebycheff coefficients (g_i) determined with (4) to (7).

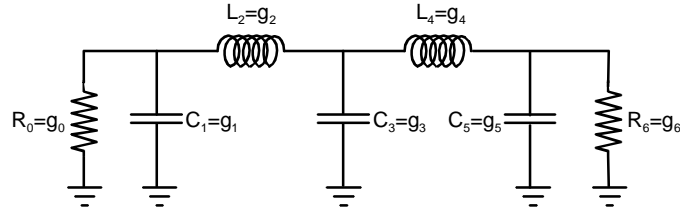


Fig. 2. 5th-order lowpass prototype filter.

$$\beta = \ln \left[\coth \left(\frac{|Las|}{17.37} \right) \right] \quad \gamma = \sinh \left(\frac{\beta}{2N} \right) \quad (4)$$

$$a_{i|i=1..N} = \sin \left(\frac{(2i-1)\pi}{2N} \right) \quad b_{i|i=1..N} = \gamma^2 + \sin^2 \left(\frac{i\pi}{N} \right) \quad (5)$$

$$g_0 = 1 \quad g_1 = \frac{2a_1}{\gamma} \quad g_{i|i=2..N} = \frac{4a_i a_{i-1}}{b_{i-1} g_{i-1}} \quad (6)$$

$$\begin{cases} g_{N+1} = 1 & \text{for } N \text{ odd} \\ g_{N+1} = \coth^2 \left(\frac{\beta}{4} \right) & \text{for } N \text{ even} \end{cases} \quad (7)$$

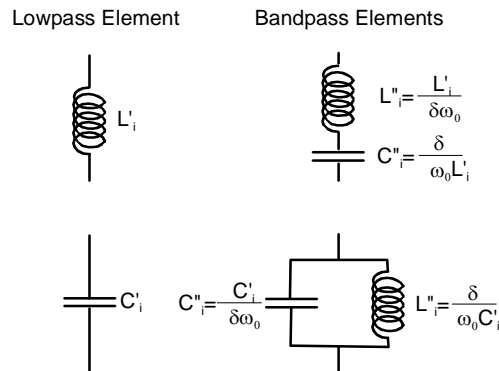


Fig. 3. Lowpass to bandpass transformation.

This lowpass network is normalized with respect to a $1\text{-}\Omega$ load and source resistances and has a 1 rad/s cut-off frequency. Hence, to obtain the desired frequency and impedance level, one needs to employ the proper transformations to scale the component values of the prototype networks (4)-(9). Indeed, a source impedance Z_0 ($Z_0=Z_{TE10}$) can be obtained by multiplying the impedance of the filter by Z_0 .

Moreover, the change of cutoff frequency from unity to ω_0 requires that we scale the frequency dependence of the filter by the factor $1/\omega_0$. As a consequence, considering a normalized values z_L and z_C , impedance and frequency scaling result in:

$$z_L = jL\omega \rightarrow Z_L = j \frac{Z_0 L \omega}{\omega_0} = jL' \omega \quad (8)$$

$$y_C = jC\omega \rightarrow Y_c = j \frac{Y_0 C \omega}{\omega_0} = jC' \omega \quad (9)$$

which shows that the new element values are given by

$$L' = \frac{Z_0 L}{\omega_0} \quad \text{and} \quad C' = \frac{C}{Z_0 \omega_0} \quad (10)$$

The bandpass filter elements can be obtained using the frequency substitution (2) for the series reactance and shunt susceptance of lowpass prototype.

$$jX_i = jL'_i \omega' \rightarrow j \frac{\omega L'_i}{\delta \omega_0} + \frac{\omega_0 L'_i}{j \delta \omega} = jL''_i \omega + \frac{1}{j C'_i \omega} \quad (11)$$

$$jB_i = jC'_i \omega' \rightarrow j \frac{\omega C'_i}{\delta \omega_0} + \frac{\omega_0 C'_i}{j \delta \omega} = jC''_i \omega + \frac{1}{j L'_i \omega} \quad (12)$$

Thus, a series inductor L'_i is transformed to a series L''_i - C''_i circuit and a shunt capacitor C'_i is transformed into shunt L''_i - C''_i circuit using rules described in Fig. 3. As a result, the desired bandpass filter prototype is depicted in Fig. 4.

To do so, each series reactance X_s and shunt susceptances B_p were replaced with their ABCD matrix, respectively

$$T = \begin{pmatrix} 1 & jX_s \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad T = \begin{pmatrix} 1 & 0 \\ jB_p & 1 \end{pmatrix} \quad (13)$$

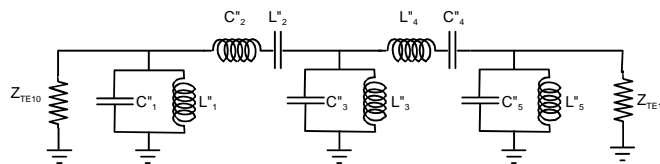


Fig. 4. 5th-order bandpass prototype filter.

TABLE I
VALUES OF THE BANDPASS FILTER LUMPED ELEMENTS

	RESONATORS 1 & 5	RESONATORS 2 & 4	RESONATOR 3
L'' (nH)	0.996	36.16	0.5794
C'' (pF)	0.2626	$6.68 \cdot 10^{-3}$	0.4522

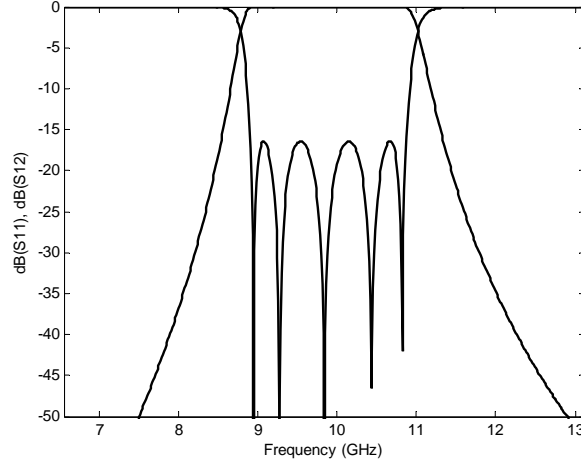


Fig. 5. Simulated results (Matlab) of a 5th-order bandpass prototype filter.

As the filter consists of cascaded series reactance and parallel susceptances, the global filter ABCD matrix is obtained by multiplying each respective matrix. One just has to transform chain matrix into scattering parameter matrix (S) as follow:

$$S = \begin{pmatrix} \frac{T_{11} + T_{12}/Z_{TE10} - T_{21} * Z_{TE10} - T_{22}}{T_{11} + T_{12}/R + T_{21} * R + T_{22}} & \frac{2 * (T_{11} * T_{22} - T_{12} * T_{21})}{T_{11} + T_{12}/Z_{TE10} + T_{21} * Z_{TE10} + T_{22}} \\ \frac{2}{T_{11} + T_{12}/Z_{TE10} + T_{21} * Z_{TE10} + T_{22}} & \frac{-T_{11} + T_{12}/Z_{TE10} - T_{21} * Z_{TE10} + T_{22}}{T_{11} + T_{12}/Z_{TE10} + T_{21} * Z_{TE10} + T_{22}} \end{pmatrix} \quad (14)$$

with R the source resistance, *i.e* Z_{TE10} mode impedance.

Each previous steps of the filter synthesis were implemented onto Matlab program. Therefore, automated calculations were ensued whatever the filter specifications. The simulation results for the studied 5th order filter are depicted in figure 5.

B. Semi-lumped Elements Design

The previous synthesis is a classical one which can be applied whatever the technology used for filter fabrication. For waveguide filters, since the passband is narrow enough, synthesis based on impedance or admittance inverters are generally used. Hence, such a technique is well supplied when an impedance level is chosen for the device implementation. In our case, as our filter is implemented in waveguide technology, to seek for simplicity, the waveguide sections were kept constant. As a consequence, the resonators which constitute the filter have constant impedance equal to Z_{TE10} .

1) Design of bandpass filters with impedance inverters

Considering the lowpass topology depicted in figure 2, such a topology can be transformed using impedance inverters. Hence, the modified prototype (Fig. 6) is composed of six impedance inverters (K_{ij}) and five lumped elements (L_i). Inverter values are calculated as follow:

$$K_{01} = Z_{TE10} \cdot \sqrt{\frac{\ell_1}{g_0 \cdot g_1}} \quad K_{56} = Z_{TE10} \cdot \sqrt{\frac{\ell_5}{g_5 \cdot g_6}} \quad (15)$$

for input and output inverters, and for other inverters

$$K_{i,i+1} \Big|_{i=2..N-1} = Z_{TE10} \sqrt{\frac{\ell_i \cdot \ell_{i+1}}{g_i \cdot g_{i+1}}} \quad (16)$$

where g_i are the Tchebycheff coefficients, the normalized inductance values ℓ_i being chosen arbitrary, in our case they were kept constant ($\ell_1 = \ell_2 = \ell_3 = \ell_4 = \ell_5 = \ell$).

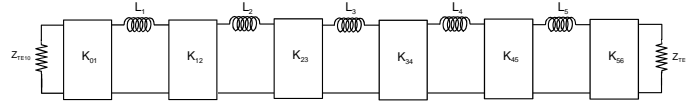


Fig. 6. 5th-order lowpass prototype filters with impedance inverters.

Transformation of lowpass prototype to bandpass prototype is made in the same way as in previous study. That is, series inductance are transformed to series L' - C' circuit (Figure 7). With regard to impedance inverters, according to (15) and (16), they are kept constant with frequency transformation (Figure 7).

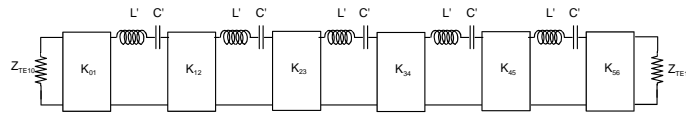


Fig. 7. 5th-order bandpass prototype filters with impedance inverters.

2) Technological implementation

The bandpass filter prototype with impedance inverters is composed of series L' - C' circuits. Due to Richard's transformations, in the vicinity of $l \approx \lambda g/2$ ($\theta = \beta l \approx 180^\circ$), the series element behave like series resonant circuit. In our concern to keep impedance levels to fundamental mode impedance, a suitable inductance value has to be chosen for the base lowpass network. So as to ensure equivalence between LC resonant circuit and half-wavelength transmission line, one has to equalize the slope parameter of their equivalent reactance. It was shown in [6] that reactance of a half wavelength transmission line can be expressed as:

$$X = Z_0 \cdot \frac{\pi}{2} \left(\frac{\omega}{\omega_0} - \frac{\omega_0}{\omega} \right) \quad (17)$$

Thus, its slope parameter is

$$S = \frac{\omega}{2} \frac{\partial X}{\partial \omega} \Big|_{\omega=\omega_0} = Z_0 \frac{\pi}{2} \quad (18)$$

Moreover, the reactance of L'C' equivalent circuit is

$$X_{eq} = \omega L' - \frac{1}{\omega C'} \quad (19)$$

which, considering resonant condition ($L'C'\omega^2=1$), results in a slope parameter of:

$$S_{eq} = \frac{\omega}{2} \frac{\partial X_{eq}}{\partial \omega} \Big|_{\omega=\omega_0} = L' \omega_0 \quad (20)$$

Hence, equivalence is obtained for inductance value for the passband prototype of:

$$L' = \frac{\pi Z_0}{2 \omega_0} \quad (21)$$

which means that, the normalized inductance value that has to be taken into account in the base lowpass prototype inverter calculation is (22), with regards to transformations depicted in figure 3:

$$\ell = \frac{\delta \pi}{2} \quad (22)$$

It only remains to transform inverter into technology. Impedance inverters can be modeled using T-network based on parallel inductance and transmission line sections (Fig. 8). The equivalence is easily verified by equating ABCD matrices of two networks. By this way, equivalence is linked with formulas (23) to (25).

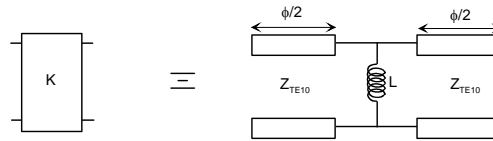


Fig. 8. Equivalent network of impedance inverter.

$$K = Z_{TE10} \tan \left| \frac{\phi}{2} \right| \quad (23)$$

$$\phi = -\arctan \frac{2X_L}{Z_{TE10}} \quad (24)$$

$$\frac{X}{Z_{TE10}} = \frac{K/Z_{TE10}}{1 - \left(K/Z_{TE10}\right)^2} \quad (25)$$

As a consequence, inverters can be realized using inductive irises and additional transmission lines. The final bandpass filter can be synthesized as depicted in Fig. 9, with effective resonators lengths given by (26):

$$L_{res_i} \Big|_{i=1..N} = \frac{\lambda_g}{2} \left(1 + \frac{\phi_{i-1,j} + \phi_{i,j+1}}{2\pi} \right) \quad (26)$$

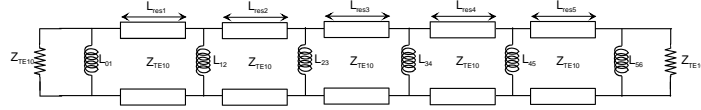


Fig. 9. 5th order Bandpass filter network.

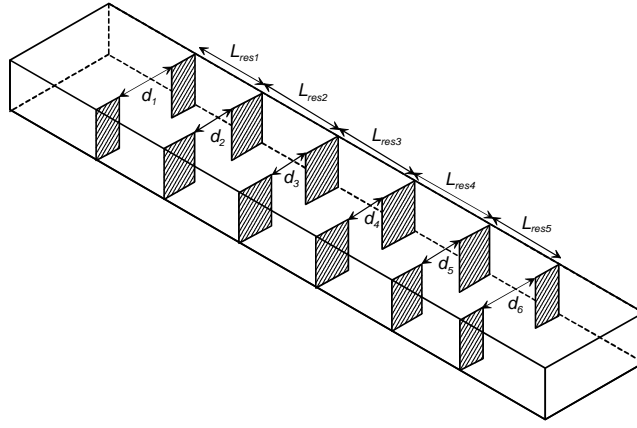


Fig. 10. 5th order Bandpass semi-lumped elements filter simulation results (Matlab).

TABLE 2

VALUES OF THE 5TH BANDPASS FILTER SEMI-LUMPED ELEMENTS

WAVEGUIDE LENGTH (mm)	SUSCEPTANCE	INDUCTANCE (nH)
$L_{res1}=L_{res5}=15.7$	$b_1=b_6=1.3872$	$L_1=L_6=4.17$
$L_{res2}=L_{res4}=17.6$	$b_2=b_5=3.7411$	$L_2=L_5=1.527$
$L_{res3}=18$	$b_3=b_4=5.0474$	$L_3=L_4=1.132$

III. MATLAB SIMULATION RESULTS

The global filter was simulated onto Matlab with nominal dimensions given in Table 2. One should note that the filter bandwidth was significantly reduced (Fig. 10). This is due to the equivalent model used for impedance inverters. Hence, such a model is usually used since the bandwidth is narrow enough. In the studied case, as the relative bandwidth to be reached is 20%,

the model is out of the model validity domain, approximation based on the slope parameter being available in the vicinity of the central frequency.

However, with regard to the slightest correction that has to be made, a simple way consists in modifying the filter element values. Indeed, as the bandwidth is reduced, the resonators are under-coupled. One just has to correct the inductive the susceptance values changes and therefore the resonator lengths. Such a method allows students tackling problems due to under- or over-coupled resonator as well as the best way to manage couplings while respecting return loss levels. The electrical response of the corrected filter is depicted in Figure 11.

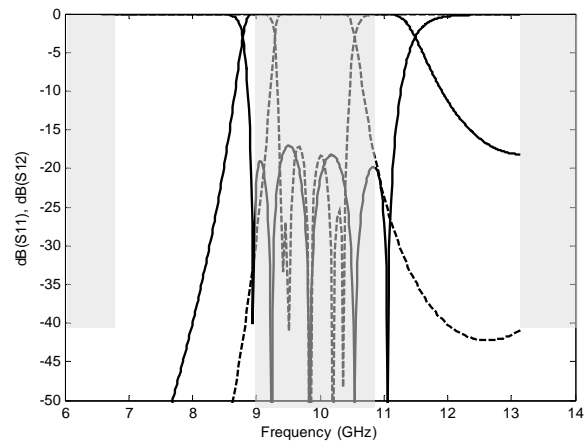


Fig. 11. 5th order Bandpass semi-lumped elements filter simulation results (Matlab). Comparison between corrected (dot line) and nominal filter.

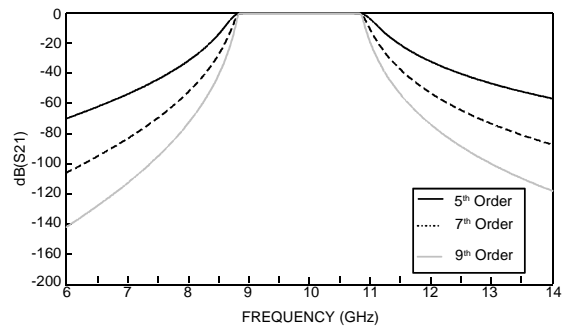


Fig. 12. Evolution electrical response vs order of ideal Tchebycheff filter.

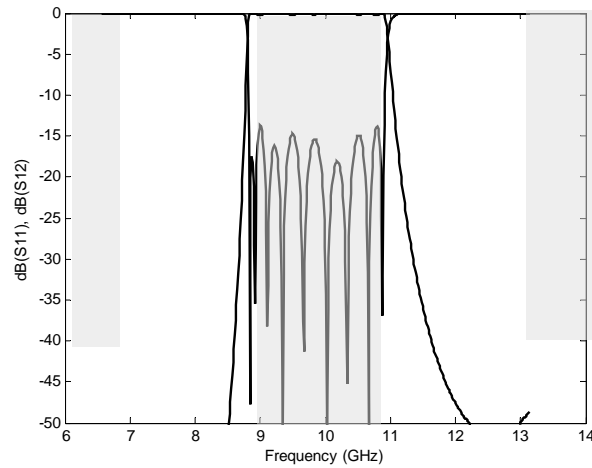


Fig. 13. 9th order Bandpass semi-lumped elements filter simulation results.

TABLE 3

VALUES OF THE 9th-ORDER BANDPASS FILTER SEMI-LUMPED ELEMENTS

WAVEGUIDE LENGTH (mm)	SUSCEPTANCE	INDUCTANCE (nH)
$L_{res1}=L_{res9}=14.8$	$b_1=b_{10}=0.9383$	$L_1=L_{10}=6.087$
$L_{res2}=L_{res8}=16.3$	$b_2=b_9=1.9913$	$L_2=L_9=2.868$
$L_{res3}=L_{res7}=16.8$	$b_3=b_8=2.5067$	$L_3=L_8=2.279$
$L_{res4}=L_{res6}=16.9$	$b_4=b_7=2.6441$	$L_4=L_7=2.160$
$L_{res5}=17$	$b_4=b_6=2.7441$	$L_4=L_7=2.081$

Even though corrected filter meets nominal specifications in the passband, it is not the same for out-of-band response. This is a classical fact when designing filter in waveguide technology since Tchebycheff theory does not take the transmission lines frequency dependence into account. Hence, due to Richard transformations, as the resonators electrical lengths are $\lambda g/2$, the filter response is periodic in frequency, repeating every f_0 . Thus, the first filter harmonic occurs at $2.f_0$, *i.e* 19.68-GHz. As the bandwidth is 20%, the out-of-band range is 10.824 GHz-17.126 GHz, which is too close to ensure 40-dB rejection at $f_i=13$ -GHz.

Thereby, one has to increase the filter order so as to ensure the desired out-of-band specification. According to the Tchebycheff filter synthesis, it is in our interest to have an odd-order filter. Indeed, with even-order filters, input and output loads are not the same. In our concern to keep waveguides with constant impedances, the filter we developed is of the 9th-order. Indeed, as depicted in Fig. 11, a 7th order filter only allows 20-dB improvement which is not sufficient enough to meet the desired specification. The design procedure is the same as previously described. The final filter simulation results are depicted in

Fig. 12, with its elements values given in Table 3.

IV. THIN INDUCTIVE DIAPHRAGM MODEL

The thin inductive diaphragm determination illustrates the modeling method delivered to the students. In this course, electromagnetic basis are taught such as waveguide principles and characteristics, modal basis decomposition, boundaries conditions, symmetry plans, trial functions in order to determine analytical formulation of chosen parameter for “simple” structure. The students must follow several steps: analyze the structure and simplify it if possible, determine an electrical equivalent circuit model which represents in a modal formulation the structure, program it and analyze the results from a physical point of view. In this example, the iris is analyzed in an infinite rectangular waveguide to qualify the diaphragm (Fig. 14). The study rests on three steps :

Step 1: considerations on this structure can be underlined to simplify the problem formulation. In the proposed frequency range, the only mode that propagates is the fundamental mode TE_{10} which admits a magnetic symmetry plan for $x=0$ and no variations along the y -axis. In the iris field, the electric field can be described by a TE_{10} mode in the iris aperture as a first approximation, it also admits a magnetic symmetry plan in $x=0$ and no variations along the y -axis. Therefore the problem does not depends on the y variable since $TM_{n,0}$ modes do not exists in rectangular waveguides. Only the $TE_{n,0}$ modes can be excited when the TE_{10} mode is scattered in the diaphragm.

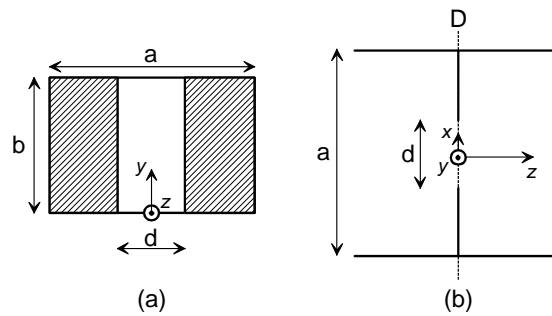


Fig. 14. Inductive diaphragm geometry (a), longitudinal view (b).

Let us consider the source and the iris field, the modes involved in the diaphragm description must present a magnetic plan in $x=0$. Hence, the only modes that should be considered are the odd modes. With magnetic symmetry in $z=0$, analyze can be limited to the half structure. Thus, the TE_{10} mode in the rectangular waveguide ($a \times b$) can be considered as a source for the problem, the field description in the iris aperture is assimilated to a TE_{10} mode in the rectangular waveguide ($d \times b$), $TE_{2p+1,0}$ modes are the only modes that are scattered in the diaphragm.

Step 2: the iris insertion in the rectangular waveguide is then represented by an equivalent circuit depicted in Fig. 15. Within this model, the TE_{10} waveguide excitation is represented by a current source J_0 (27), the modal operator \hat{Y} (28) represents

$TE_{2p+1,0}$ evanescent modes excited by TE_{10} scattered around the iris and the field in the iris is represented by a virtual source E_e (29) which is called virtual because no power is delivered to the waveguide.

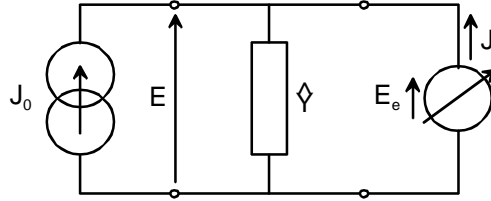


Fig. 15. Boundary condition representation by equivalent circuit for iris aperture in rectangular waveguide.

$$J_0 = I_0 f_0 \quad (27)$$

$$\hat{Y} = \sum_{p=1}^{+\infty} |f_p\rangle Y_{2p+1,0}^{TE} \langle f_p| \quad (28)$$

$$E_e = v_e g_e = v_e \cos\left(\frac{\pi x}{d}\right) \vec{y} \quad \text{for } x \in \left[-\frac{d}{2}, \frac{d}{2}\right] \quad (29)$$

with f_p the modal basis function of the rectangular waveguide (30), $Y_{2p+1,0}^{TE}$ the modal admittance of the evanescent modes (31), v_e the field amplitude in the iris and I_0 the source amplitude.

$$f_p = \sqrt{\frac{\tau_p}{a}} \cos\left(\frac{(2p+1)x}{a}\right) \vec{y} \quad \text{for } x \in \left[-\frac{a}{2}, \frac{a}{2}\right] \quad (30)$$

$$Y_{2p+1,0}^{TE} = \frac{\sqrt{\left(\frac{2p+1}{a}\pi\right)^2 - k_0^2}}{j\omega\mu_0} \quad (31)$$

with $\tau_p = 2$ if $p \neq 0$ else $\tau_0 = 1$.

Step 3: the circuit resolution leads to the system (32) which is easily solved into (33).

$$\begin{bmatrix} E \\ J \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & \hat{Y} \end{bmatrix} \begin{bmatrix} J_0 \\ E_e \end{bmatrix} \quad (32)$$

$$[Z] = \frac{Z_{TE}}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \quad (33)$$

with V_0 the TE_{10} (source) amplitude; Z_{TE} (34) the input impedance deduced from the equivalent circuit.

$$Z_{TE} = \frac{V_0}{I_0} = \frac{|f_0\rangle g_e|^2}{\sum_{p=1}^{+\infty} Y_{2p+1,0}^{TE} |f_p\rangle g_e|^2} = jL_{TE}\omega \quad (34)$$

This method based on equivalent circuits is easily accessible by the students because it is a systematic and rigorous procedure which presents the different steps of the electromagnetic approach by a simple scheme.

Step 4: The L_{TE} parameter is programmed on Matlab and compared to the first approximation usually used (35) which is obtained with conventional mode matching technique [7].

$$L = \frac{1}{bY_{1,0}^{TE}} \quad (35)$$

with the susceptance b defined by

$$b = \frac{\lambda_g}{a} \cot^2\left(\frac{\pi d}{2a}\right) \quad (36)$$

Both methods were simulated and results were compared in Fig. 16. The error represented in Fig. 17 is important whatever the iris dimension. A first order approximation is not sufficient to describe well the value of inductive diaphragm. As we will see in the next part, the use of such an approximation implies significant correction in the filter design.

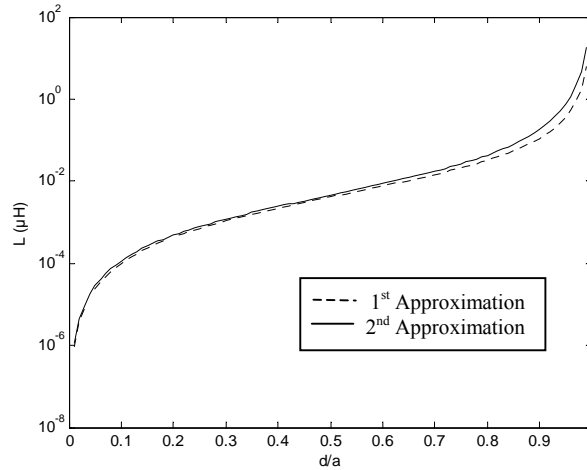


Fig. 16. Comparison between simulated susceptance value with classical mode matching technique and electromagnetic-based equivalent circuit.

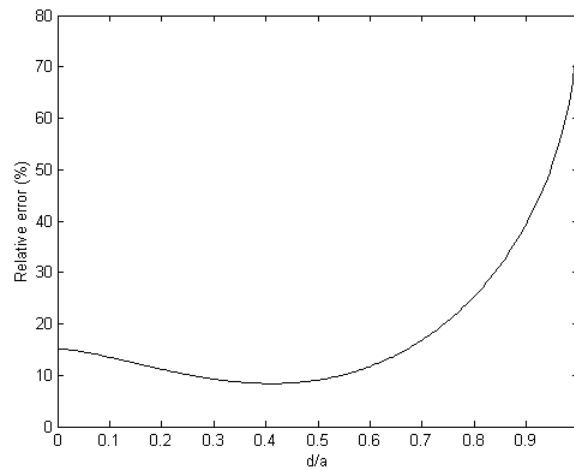


Fig. 17. Relative error between the two methods inductance evaluation.

V. ELECTROMAGNETIC SIMULATIONS AND MEASUREMENTS

Based on the previous study, the dimensions of iris apertures were evaluated with both classical model and our model based on modal analyses (Table 4). An electromagnetic simulation is performed with Ansoft-HFSS. Such a computation method is necessary since circuit analyses did not take the higher order modes into account, those modes being generated by the iris discontinuities. One can observe a good agreement between the expected specifications for the filter designed with optimized model. Indeed, both passband and rejection frequency are met. This is not the same for filter designed with conventional mode matching model. Hence, out-of-band specification is met, nevertheless, the passband is narrow and the lower cutoff frequency is not ensured. Thus, the iris aperture model we developed should henceforth be used for future filter design.

Once the filter dimensions were obtained, the filter was realized and measured. Experimental results showed a slight frequency shift below 2%, as well as a return loss debasement in the beginning of the passband. This can easily be explained by the inaccuracy of the technological process used.

TABLE 4

DIMENSIONS OF THE 9TH-ORDER BANDPASS FILTER

WAVEGUIDE LENGTH (mm)	IRIS APERTURE (mm) MODEL I	IRIS APERTURE (mm) MODEL II
$L_{\text{res}1}=L_{\text{res}9}=14.8$	$d_1=d_{10}=13.7391$	$d_1=d_{10}=14.197$
$L_{\text{res}2}=L_{\text{res}8}=16.3$	$d_2=d_9=11.0414$	$d_2=d_9=11.385$
$L_{\text{res}3}=L_{\text{res}7}=16.8$	$d_3=d_8=10.2097$	$d_3=d_8=10.533$
$L_{\text{res}4}=L_{\text{res}6}=16.9$	$d_4=d_7=10.0187$	$d_4=d_7=10.338$
$L_{\text{res}5}=17$	$d_5=d_6=9.8864$	$d_5=d_7=10.2035$

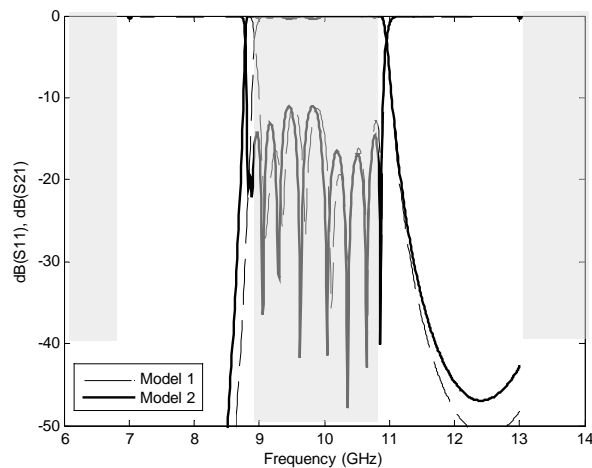


Fig. 18. Electromagnetic simulation results (HFSS). Comparison between conventional mode matching and electromagnetic-based equivalent circuit techniques.

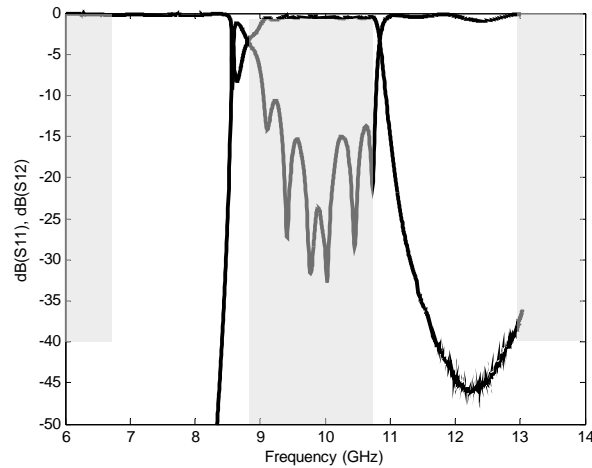


Fig. 19. Measurement results of the 9th-order bandpass filter.

Hence, irises were manually realized and welded in the wave guide. Thus, the ensued technological dispersion is of the order of 1% which is not suited with the global filter sensitivity and complexity. This was shown with post-simulation performed with Ansoft-HFSS.

VI. CONCLUSION

A global design procedure was developed for bandpass filter implemented in waveguide technology with inductive iris coupling aperture. Based on classical synthesis, technological implementation was investigated which highlighted difficulties met in the design of devices in microwave frequency range. Discontinuities model were developed using electromagnetic equivalent-circuit method which accuracy was proved in comparison with 3D electromagnetic simulation as well as in experiment.

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